# UV radiation in Ar-O<sub>2</sub>, N<sub>2</sub>-O<sub>2</sub> and Ar-O<sub>2</sub>-N<sub>2</sub> microwave discharges and post-discharges

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This work presents a theoretical investigation on the production of active species in low-pressure  $Ar-O_2$  and  $Ar-O_2-N_2$  microwave discharges, and the evolution of the species densities along the afterglow. Special attention is given to the UV emitting species, such as Ar(4s) excited states, and NO(A) and NO(B) molecules. It is shown that, for the conditions of the present study, the active species are efficiently produced in the discharge, however their lifetime in the afterglow is limited. We have shown that the Ar(4s) excited states depopulate during  $10^{-4}$  s afterglow time in pure Ar, and the depopulation becomes faster with  $O_2$  addition. In what concerns the NO(A) and NO(B) molecules, their densities increase with the  $N_2$  addition to the Ar-2%O<sub>2</sub> mixture. The NO(A) density has found to be maximum in the discharge in 98%N2-2%O2 mixture, as well as the N and O atoms densities.

# 1. Introduction

Post-discharge plasmas have a wide range of applications, such as metal surface cleaning, medical sterilization, etching and grafting, thin film synthesis, etc. The post-discharges of Ar-O<sub>2</sub> and N<sub>2</sub>-O<sub>2</sub> flowing microwave discharges have been found to be very efficient in the bacterial spore removal [1-3]. Kylián et al. [4] have studied a low pressure Ar-O<sub>2</sub>-N<sub>2</sub> ICP discharge in order to evaluate its properties in terms of sterilization. They have demonstrated that the application of Ar-O<sub>2</sub>-N<sub>2</sub> ternary mixture offers the possibility to combine the advantageous properties of the above-mentioned binary mixtures. The main sterilizing agents in Ar- $O_2$  and  $N_2$ - $O_2$  plasmas it is believed [1,3,5] to be the O atoms and UV photons. The aim of the present study is to investigate the evolution of the UV radiation when changing from binary mixtures to the ternary mixture.

#### 2. System set-up

The post-discharge system used in our investigations is similar to the one used in the sterilization experiments by the Montréal group [6]. Here the surface wave discharge is generated in a R=5 mm silica tube with 2450 MHz frequency field. The discharge and the reactor are linked through two tubes of different diameters, whereas the first tube, right after the discharge zone, has the same diameter as the discharge tube, the second tube

is widened to 2.6 cm relatively to the discharge tube. Herein we focus our attention on the discharge and early afterglow region and assume that the radius of the tube where the early-afterglow develops has the same size as the discharge tube. Our calculations are conducted in Ar-O<sub>2</sub>, N<sub>2</sub>-O<sub>2</sub> and Ar-O<sub>2</sub>-N<sub>2</sub> mixtures, respectively, for different mixture compositions at 2 Torr pressure.

# 3. Model

The system is described with two different models valid for the discharge and early afterglow region, and for the late afterglow present in the large reactor, respectively. The species densities in the discharge region are calculated by solving the homogeneous electron Boltzmann equation, coupled together with the rate balance equations of different species. The concentrations obtained for the steadystate discharge are used as initial values to the early afterglow taking place in the tube connecting the discharge to the main reactor, where the same system of equations is solved in time under zero electric filed [7,8]. The evolution of the species densities in the post-discharge reactor is followed with a 3-D hydrodynamic model [9].

In the model of the Ar-O<sub>2</sub> system the species taken into account are: Ar( ${}^{1}S_{0}, {}^{3}P_{2}, {}^{3}P_{1}, {}^{3}P_{0}, {}^{1}P_{1}$ ), O<sub>2</sub>(X<sup>3</sup> $\Sigma_{g}$ ,v), O<sub>2</sub>(a<sup>1</sup> $\Delta_{g}, b^{1}\Sigma_{g}^{+}$ ), O( ${}^{3}P, {}^{1}D$ ), O<sub>3</sub>, Ar<sup>+</sup>, Ar<sub>2</sub><sup>+</sup>, O<sub>2</sub><sup>+</sup>, O<sup>+</sup>, O<sup>-</sup>; while in the Ar-N<sub>2</sub>-O<sub>2</sub> system to the above listed species the N containing species are still added, i.e.: N<sub>2</sub>(X,v), N({}^{4}S, {}^{2}D, {}^{2}P),

 $N_2(A,B,B',C,a',a,w),$  $NO_2(X, A),$ NO(X, A, B).  $N_2^+(X, B)$ ,  $N_4^+$  and NO<sup>+</sup>. In the hydrodynamic model of the late-afterglow we neglect some of the excited species and all the ions, since due to deactivation and recombination their densities become negligible in the reactor. The set of gas phase reactions describing the kinetics of species as considered in the post-discharge chamber and discharge region for the nitrogen and oxygen containing species has been presented in [8,10] and the references therein. A set of reactions for the Ar, Ar-N<sub>2</sub> and Ar-O<sub>2</sub> mixtures can be found in [12-14]. The surface reactions involving the N and O atoms have been discussed in details in [11].

#### 4. Results

The UV radiation in Ar-O<sub>2</sub> discharge is due to the  $Ar({}^{3}P_{1}, {}^{1}P_{1})$  resonant states, while in N<sub>2</sub>-O<sub>2</sub> mostly to the NO(A) and NO(B) molecules.

As a first step we investigate the densities of  $Ar({}^{3}P_{1}, {}^{1}P_{1})$  resonant states in Ar-O<sub>2</sub> mixture, and the evolution of these densities in the early-afterglow region. Figure 1 shows the calculated  $Ar({}^{3}P_{1}, {}^{1}P_{1})$ densities as a function of the early-afterglow time (flight-time of species in the afterglow, where t=0 means the end of the discharge) in Ar and in Ar-O<sub>2</sub> mixture with 1% and  $2\%O_2$  addition. The calculations show that the in pure Ar the resonant states depopulate within 0.1 ms in the earlyafterglow. Moreover, with oxygen addition the depopulation occurs faster. This is a consequence of to the quenching of Ar 4s  $({}^{3}P_{2}, {}^{3}P_{1}, {}^{3}P_{0}, {}^{1}P_{1})$  states by through reaction  $Ar(4s)+O_2 \rightarrow Ar(^{1}S_0)+O(^{3}P)$  $O_2$  $+O(^{3}P, ^{1}D)$ . The results here obtained suggest that in the late afterglow present in a large reactor no UV emission can be expected from the  $Ar({}^{3}P_{1}, {}^{1}P_{1})$ states.

Figure 2 shows the evolution of oxygen  $O({}^{3}P)$  and  $O({}^{1}D)$  atoms in the afterglow. We can notice the steep increase of the O atoms upon oxygen addition into an argon discharge, and that oxygen is strongly dissociated by electron impact and by the Ar metastables (above mentioned reaction), the dissociation degree being in the range of 25-50%. Similar results have also been obtained experimentally by Mozetič et al. [15] where the density of  $O({}^{3}P)$  ground state atoms has been determined both by NO titration and catalytic probe methods.



Figure 1. Time evolution of the Ar atoms resonant states densities along the early afterglow (t=0 means the end of the discharge) in Ar (thick line) and in Ar-O<sub>2</sub> mixture with 1% and 2%O<sub>2</sub> addition (thinnest line).



Figure 2. Time evolution of  $O({}^{3}P)$  (thick lines) and  $O({}^{1}D)$  (thin lines) atoms densities along the early afterglow in Ar-O<sub>2</sub> mixtures with different O<sub>2</sub> addition.

As a second step N<sub>2</sub> (from 0.2 to 98%) was added into the Ar-O<sub>2</sub> mixture while the O<sub>2</sub> percentage in the mixture was kept constant at 2%. Figure 3 a) and b) shows the evolution of NO(A) and NO(B) densities in the early-afterglow. In the discharge the NO(A) molecules are mainly created from the collision of NO(X) with N<sub>2</sub>(A) metastable molecules: NO(X)+N<sub>2</sub>(A) $\rightarrow$ NO(A)+N<sub>2</sub>(X), while NO(B) molecules are created through the three body re-association of N and O atoms: N+O+M  $\rightarrow$ NO(B)+M [6]. Consequently, in the discharge NO(A) is always significantly more populated than NO(B) and the time evolution of these two states is quite different in the early afterglow, up to times 10<sup>4</sup>-10<sup>-3</sup> s, as illustrated in Figure 3.



Figure 3. Time evolution of (a) NO(A) and (b) NO(B) molecules densities along the early afterglow in  $Ar-O_2-N_2$  mixtures with different  $N_2\%$  addition into  $Ar-2\%O_2$ .

The calculations show that the highest NO(A) density is obtained in the  $98\%N_2-2\%O_2$  mixture. However at around 0.1 ms in the afterglow this density decreases below the one obtained in mixtures of 50% and 90% N<sub>2</sub>. From this position in the afterglow NO(A) molecules are dominantly created through the three body re-association of N and O atoms, similarly to the case of NO(B) molecules [6]. Therefore, the densities of NO(A) and NO(B) UV emitting molecules in the afterglow are controlled by the N and O atoms densities and, for longer afterglow times, have a rather similar evolution. Notice that N<sub>2</sub> concentrations and higher strongly favour the presence of NO(A) molecules in the late afterglow.

Figure 4 a) and b) shows the O and N atoms densities, respectively, along the ealy-afterglow for different  $N_2$  addition to the Ar-2%O<sub>2</sub> mixture. The results show that the atomic species N and O do not change significantly in the early afterglow, and both increase with  $N_2$  addition into the mixture.

As the oxygen concentration is fixed to 2%, the rate of O atom production in the discharge by electron



Figure 4. Time evolution of (a)  $O(^{3}P)$  and (b)  $N(^{4}S)$  atoms densities along the early afterglow in Ar-O<sub>2</sub>-N<sub>2</sub> mixtures with different N<sub>2</sub>% addition into Ar-2%O<sub>2</sub>.

impact is nearly constant when nitrogen is added into the mixture. Near pure Argon, an additional atomic oxygen source is dissociation due to collisions with the Ar(4s) states, as referred above. Evidently, this mechanism looses importance for higher N<sub>2</sub> concentrations in the mixture. However, the disappearance of this process is compensated by the raise in the contribution of other dissociation mechanisms, which become influential for significant nitrogen percentages, such as  $O_2$ dissociation in collisions with N<sub>2</sub> excited molecules,  $N_2(B, a, a')+O_2 \rightarrow N_2+O+O$ . The dominant reaction is the one involving  $N_2(B)$  molecules. As an outcome, higher O atom concentrations are observed for higher nitrogen percentages in the mixture, as it can be seen in figure 4 a).

In turn, the N atom density increase in the discharge with  $N_2$  content revealed by figure 4 b) is basically justified by the direct dissociation of  $N_2$  molecules, the density of the latter course increasing. Nitrogen molecules are mostly dissociated by direct electron impact and, to a smaller extent, by collisions between  $N_2(A)$ 

molecules and vibrationally excited molecules.  $N_2(X,v\geq 13)+O\rightarrow NO+N$ , Reaction although exhibiting a very absolute high rate, does not contribute to the creation of N atoms as it does not prevail process over the reverse  $N+NO \rightarrow N2(X,v\approx 3)+O$ [16]. Another relevant dissociation mechanism is the destruction of NO via  $N_2(a')+NO \rightarrow N_2+N+O.$ 

The present investigation casts light into the elementary processes responsible for the creation of the UV emitting species in the afterglow of the Ar- $O_2$ - $N_2$  plasmas used in plasma sterilization reactors. It is shown the kinetics of these species is strongly coupled with the atomic species N and O. For the moment we have addressed the characterization of the discharge and early afterglow. The study of the species distribution in an actual reactor will be performed in the near future.

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